

# Positive-Bias Temperature Instability (PBTI) of GaN MOSFETs

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**Abstract** — We have investigated the stability of the gate stack of GaN n-MOSFETs under positive gate stress. Devices with a gate dielectric that consists of pure SiO<sub>2</sub> or a composite SiO<sub>2</sub>/Al<sub>2</sub>O<sub>3</sub> bilayer were studied. Our research has targeted the evolution of threshold voltage ( $V_T$ ), subthreshold swing (S) and transconductance ( $g_m$ ) after positive gate voltage stress of different duration at different voltages and temperatures. We have also examined the recovery process after the stress is removed. We have observed positive  $V_T$  shift ( $\Delta V_T$ ) in both gate dielectrics under positive gate stress. In devices with a SiO<sub>2</sub> gate oxide, we have found that  $\Delta V_T$  is caused by a combination of electron trapping in pre-existing oxide traps and interface trap generation. In devices with a composite SiO<sub>2</sub>/Al<sub>2</sub>O<sub>3</sub> gate oxide, on the other hand,  $\Delta V_T$  is due to electron trapping in pre-existing oxide traps and generation of near interface oxide traps.

**Keywords**-component; GaN; MOSFETs; BTI;  $V_T$  shift

## I. INTRODUCTION

GaN power transistors represent a promising alternative to Si devices for power switching applications [1]. Compared to other semiconductor systems, GaN-based materials offer superior characteristics and performance under high-frequency and high-temperature conditions [2-4]. Recently, GaN-high electron mobility transistors with insulated-gate (MIS-HEMTs) have attracted much attention because they offer high current, high breakdown voltage and low gate leakage current, all desirable attributes for power transistors [5-7]. On the other hand, GaN MIS-HEMTs face significant reliability and stability challenges. Among them, instability of the electrical characteristics, in particular  $V_T$ , after prolonged high voltage stress at high temperature is a serious problem [8-11].

The problem of  $V_T$  instability is not unique to GaN transistors. This phenomenon has been intensively studied in Si, SiC and various III-V MOS systems where it is sometimes referred to as bias-temperature instability (BTI) [12-18]. In Si n-channel MOSFETs with high-k gate dielectric, Positive Bias Temperature Instability (PBTI) is known to be caused by electron trapping in pre-existing oxide traps that causes a positive  $V_T$  shift [12, 13]. Negative-bias temperature instability (NBTI) in p-channel MOSFETs is related to interface state generation and hole trapping in the gate dielectric [14]. In SiC MOSFETs, periodic  $V_T$  shifts under alternating positive and negative gate stress have been observed and attributed to electron tunneling in and out of near-interface oxide traps [15]. InGaAs n-MOSFETs with a high-k dielectric also have shown positive  $V_T$  shifts after positive gate stress. This is believed to

be due to pre-existing oxide traps and electron trap generation near the oxide/InGaAs interface [16-18].

For GaN MIS-HEMTs, a few studies have reported  $V_T$  shifts due to PBTI [8-10]. However, its origin is not well understood. This is partly due to the complicated nature of the gate stack of a MIS-HEMT which contains several interfaces.

To better understand the mechanisms responsible for  $V_T$  shift under PBTI, a GaN MOSFET structure is desirable because of its single oxide/GaN interface. In [11], the  $V_T$  instability under PBTI in GaN MOSFETs with a SiN gate dielectric was studied. This work revealed the importance of the conduction band discontinuity at the oxide-semiconductor interface. However, the study did not include an evaluation of the changes in the subthreshold swing nor the transconductance of the device. These factors are important in themselves and greatly illuminate the relevant physics that affect the stability of  $V_T$ .

In this work, we examine the dynamics of  $V_T$ , S and  $g_m$  in GaN MOSFET before, during and after stress. This allows us to isolate and identify different mechanisms for  $V_T$  shift. We examine two different gate dielectrics: SiO<sub>2</sub> and a SiO<sub>2</sub>/Al<sub>2</sub>O<sub>3</sub> composite.

## II. EXPERIMENTAL

### A. Device structure

The devices studied in this work are sketched in Fig. 1 (not to scale). This is a simple recessed-gate structure with an intrinsic gate stack that consists of either 50 nm of SiO<sub>2</sub> or 40 nm EOT composite SiO<sub>2</sub>/Al<sub>2</sub>O<sub>3</sub> dielectric (Al<sub>2</sub>O<sub>3</sub> next to semiconductor). The devices have channel width/length of 100  $\mu\text{m}/1 \mu\text{m}$ . The standard figures of merit in the saturation regime ( $V_{DS} = 10 \text{ V}$ ) are for SiO<sub>2</sub> devices:  $V_T \sim -1.0 \text{ V}$ , S  $\sim 175 \text{ mV/dec}$ ,  $g_{m, \text{max}} \sim 20 \text{ mS/mm}$ ; for SiO<sub>2</sub>/Al<sub>2</sub>O<sub>3</sub> devices:  $V_T \sim -0.7 \text{ V}$ , S  $\sim 170 \text{ mV/dec}$ ,  $g_{m, \text{max}} \sim 30 \text{ mS/mm}$ .

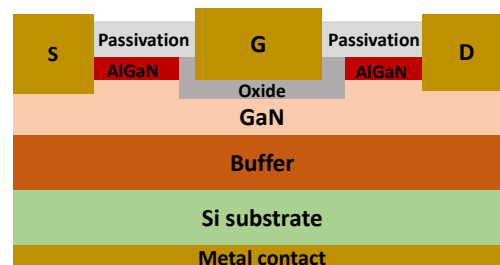


Figure 1. GaN MOSFET structure studied in this work.

In this study we focus our interest on the stability of the threshold voltage ( $V_T$ , defined at drain current  $I_D = 1 \mu\text{A}/\text{mm}$ ), the subthreshold swing ( $S$ , defined at  $I_D = 0.1 \mu\text{A}/\text{mm}$ ) and the maximum transconductance ( $g_{m,\text{max}}$ ) in the linear region (defined at  $V_{DS} = 0.1 \text{ V}$ ). Virgin devices used in this work exhibit values of  $V_T = 0.06 \text{ V}$  with a standard deviation of  $0.096 \text{ V}$ ,  $S = 110 \text{ mV}/\text{dec}$ ,  $g_{m,\text{max}} = 0.65 \text{ mS}/\text{mm}$  for  $\text{SiO}_2$  MOSFETs and  $V_T = 0.16 \text{ V}$  with a standard deviation of  $0.06 \text{ V}$ ,  $S = 100 \text{ mV}/\text{dec}$ ,  $g_{m,\text{max}} = 0.56 \text{ mS}/\text{mm}$  for  $\text{SiO}_2/\text{Al}_2\text{O}_3$  MOSFETs. Fig. 2 shows the subthreshold and transconductance characteristics at different temperatures of  $\text{SiO}_2$  and  $\text{SiO}_2/\text{Al}_2\text{O}_3$  transistors in the linear regime ( $V_{DS} = 0.1 \text{ V}$ ).

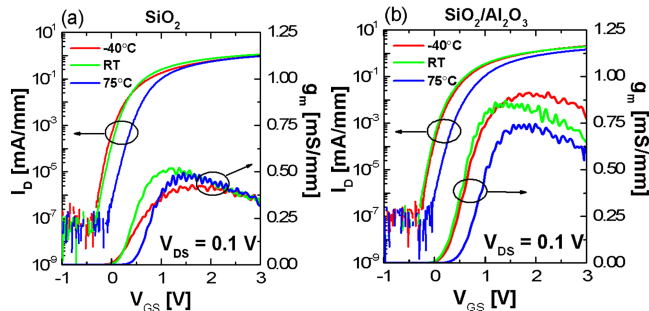


Figure 2. Subthreshold and transconductance characteristics of (a)  $\text{SiO}_2$  and (b)  $\text{SiO}_2/\text{Al}_2\text{O}_3$  GaN MOSFETs in the linear regime ( $V_{DS} = 0.1 \text{ V}$ ) at different temperatures.

### B. Experiment flow

In our studies, we first developed a benign characterization scheme that is used to characterize the devices before, during and after stress experiments. It consists of an  $I_D$ - $V_{GS}$  sweep at  $V_{DS} = 0.1 \text{ V}$  that starts at  $V_{GS} = -0.5 \text{ V}$  and stops at  $I_D = 1 \mu\text{A}/\text{mm}$ . From this we extract  $V_T$  and the subthreshold swing ( $S$ ), as defined above. This method is benign as proven by minor changes in device characteristics after 100 sweeps ( $\Delta V_T < 30 \text{ mV}$ ,  $\Delta S < 20 \text{ mV}/\text{dec}$ ).

In a typical experiment, a device is first “initialized” and then characterized. The device initialization process consists of flushing the device for 5 minutes under microscope light, followed by thermal detrapping (we use a single set of thermal detrapping conditions throughout this work). Initialization typically results in a small change in device characteristics ( $\Delta V_T < 30 \text{ mV}$ ). This step helps us create a “stable” and reproducible initial state for the device which we use as a reference for subsequent stress/recovery experiments. After initialization, we characterize the device with a complete  $I_D$ - $V_{GS}$  sweep at  $V_{DS} = 0.1 \text{ V}$  that starts at  $V_{GS} = -0.5 \text{ V}$  and stops at  $V_{GS} = 3 \text{ V}$ . From this sweep we extract the initial  $V_T$ ,  $S$  and  $g_{m,\text{max}}$ . Next, the thermal detrapping step is repeated. We then re-measure the device with a shorter, benign  $I_D$ - $V_{GS}$  sweep as described in the previous paragraph to confirm that the device characteristics are stable. At this point, the stress experiments can start.

The PBTI stress phase consists of a series of stress segments of increasing length,  $t_{\text{stress}}$ , during which the device is subjected to a constant positive gate stress ( $V_{GS,\text{stress}}$ ), with the source, drain and substrate grounded.  $V_{GS,\text{stress}}$  between 5 and

15 V and  $t_{\text{stress}}$  between 10 and 10,000 seconds have been studied at  $-40^\circ\text{C}$ , room temperature (RT) and  $75^\circ\text{C}$ . Immediately after each stress segment, the evolution of  $V_T$  and  $S$  are tracked through repeated  $I_D$ - $V_{GS}$  sweeps that are performed for  $\sim 1,000 \text{ sec}$ . The first reading comes 1 to 2 sec after the stress is stopped. At the end of each stress segment, the device is reinitialized using thermal detrapping and a complete characterization is performed. To confirm that the thermal detrapping is effective and complete, in several cases we have periodically remeasured a stressed and thermally detrapped device after storage at room temperature of various lengths of time. We have not observed any significant additional recovery of  $V_T$  within  $10^7 \text{ s}$ . We then consider the residual  $\Delta V_T$ ,  $\Delta S$  and  $\Delta g_{m,\text{max}}$  after thermal detrapping permanent. In a typical sequence, the same device is used starting from the lowest  $V_{GS,\text{stress}}$  and shortest  $t_{\text{stress}}$ . If permanent damage occurs after a stress step, i.e.,  $V_T$  and  $S$  are not restored to their initial values after thermal detrapping, we switch to a new device with well-matched initial characteristics.

In a separate set of experiments, we use a fresh collection of devices to study the behavior of  $g_{m,\text{max}}$ . We perform similar stress and recovery experiments but immediately after each stress segment we perform a one-time, downward sweep of  $I_D$ - $V_{GS}$  in the linear region ( $V_{DS} = 0.1 \text{ V}$ ) that starts at  $V_{GS} = 5 \text{ V}$  and stops at  $V_{GS} = -0.5 \text{ V}$ . Because  $V_{GS,\text{stress}}$  is at least 5 V, this sweep should not introduce additional damage to the device. Since the maximum transconductance point tends to occur between  $V_{GS} = 2.5$  and 4 V, this sweep yields a measurement of  $g_{m,\text{max}}$  in the linear regime.

## III. RESULTS

### A. Positive gate stress at room temperature (RT)

Fig. 3 (a) shows the stress time evolution of  $\Delta V_T$  under positive gate stress for  $\text{SiO}_2$  and  $\text{SiO}_2/\text{Al}_2\text{O}_3$  GaN MOSFETs at RT. Data points are extracted 1 sec after the stress is stopped. For both dielectrics, we observe positive  $\Delta V_T$  that increases with  $t_{\text{stress}}$  and  $V_{GS,\text{stress}}$ . The composite dielectric generally shows more  $V_T$  shift than the pure  $\text{SiO}_2$  gate. The evolution of  $\Delta S$  is shown in Fig. 3 (b). For the  $\text{SiO}_2$  devices, we see minimal change in  $S$  for  $V_{GS,\text{stress}} = 5$  and 10 V, but a clear increase for  $V_{GS,\text{stress}} = 15 \text{ V}$ . For the  $\text{SiO}_2/\text{Al}_2\text{O}_3$  device, we see minimal change in  $S$  at all  $V_{GS,\text{stress}}$ .

Fig. 4 shows the dynamics of  $V_T$  and  $S$  recovery at RT after  $V_{GS,\text{stress}} = 5 \text{ V}$  and  $15 \text{ V}$  for  $t_{\text{stress}} = 10,000 \text{ s}$ . We observe partial  $\Delta V_T$  recovery and no  $\Delta S$  recovery

Introducing a final detrapping step following the recovery phase allows us to assess the introduction of permanent damage during the stress phase. The open symbols in Figs. 3 and 4 show final values of  $\Delta V_T$  and  $\Delta S$  after 10,000 sec stress segments followed by thermal detrapping. The devices experience a nearly complete recovery after 5 V stress, but only partial recovery after 15 V stress. Notably,  $\Delta S$  after 15 V stress did not recover at all in the  $\text{SiO}_2$  MOSFETs.  $I_D$ - $V_{GS}$  curves before stress and after stress and thermal detrapping for  $V_{GS,\text{stress}} = 15 \text{ V}$  are shown in Fig. 5.

In general, in all studied devices, we find no significant permanent  $V_T$  shift after  $V_{GS,\text{stress}} \leq 10 \text{ V}$ , but a permanent component for  $V_{GS,\text{stress}} = 15 \text{ V}$ . For the  $\text{SiO}_2/\text{Al}_2\text{O}_3$  devices,

we observe no significant permanent S degradation. For SiO<sub>2</sub> devices, S degrades after 15 V stress and is not recoverable. This also suggests the introduction of permanent damage to these devices.

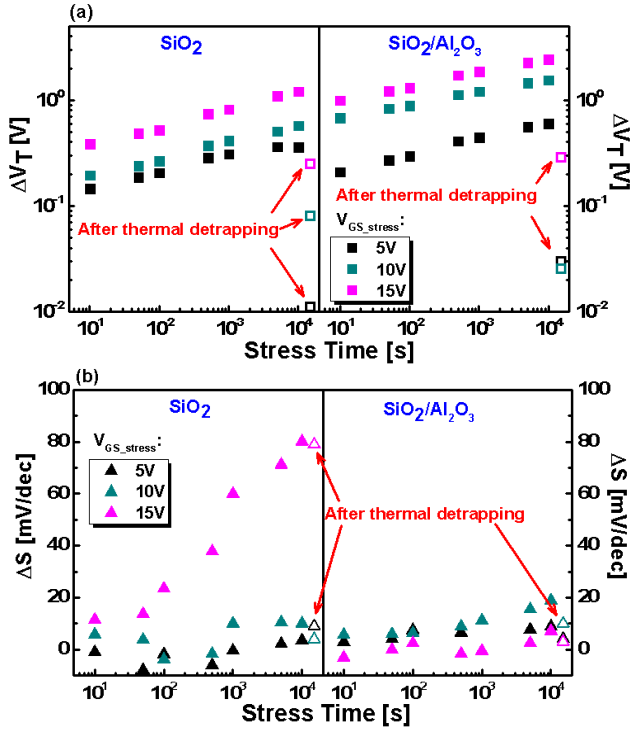


Figure 3. (a) Stress time evolution of  $\Delta V_T$ . (b) stress time evolution of  $\Delta S$ , for SiO<sub>2</sub> vs. SiO<sub>2</sub>/Al<sub>2</sub>O<sub>3</sub> GaN MOSFETs stressed at three different voltages at room temperature. Data points are taken 1 sec after the stress. The last set of points (open symbols) are  $\Delta V_T$  and  $\Delta S$  after thermal detrapping following a stress experiment with  $t_{\text{stress}} = 10,000$  s.

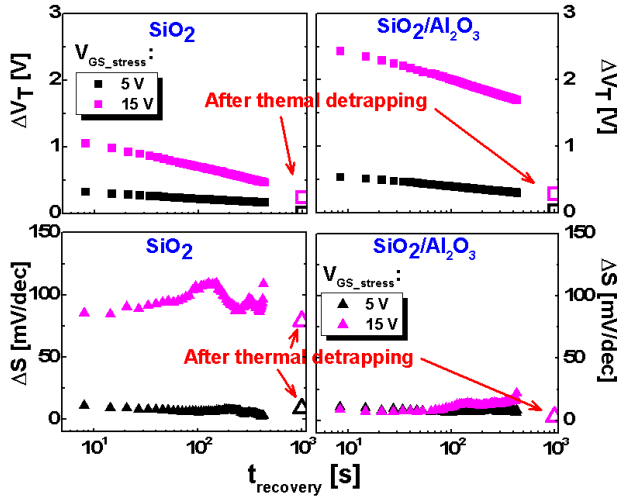


Figure 4. Recovery of  $\Delta V_T$  and  $\Delta S$  after a 10,000 sec stress at RT for SiO<sub>2</sub> vs. SiO<sub>2</sub>/Al<sub>2</sub>O<sub>3</sub> GaN MOSFETs at  $V_{GS, \text{stress}} = 5$  and 15 V. The last set of points (open symbols) are  $\Delta V_T$  and  $\Delta S$  after thermal detrapping following the recovery phase.

During each stress segment, we also monitor the evolution of the gate current ( $I_G$ ). Fig. 6 shows  $I_G$  as a function of stress

time for  $t_{\text{stress}} = 10,000$  sec experiments.  $I_G$  tends to slightly decrease with stress time. This is consistent with other studies on charge trapping in pre-existing oxide traps. It also indicates that no new trap generation is occurring [19].

Fig. 7 shows the evolution of  $g_{m, \text{max}}$  degradation ( $\Delta g_{m, \text{max}}/g_{m, \text{max}0} \sim 1$  sec after stress (closed symbols) and after subsequent thermal detrapping (open symbols) in a separate set of experiments from those of Figs. 2-6. For both dielectrics, we see higher  $g_{m, \text{max}}$  degradation with higher  $V_{GS, \text{stress}}$  and longer  $t_{\text{stress}}$ . At  $V_{GS, \text{stress}} = 5$  V,  $g_m$  degradation is minimal. After thermal detrapping,  $g_m$  completely recovers in SiO<sub>2</sub> MOSFETs, and only partially recovers in SiO<sub>2</sub>/Al<sub>2</sub>O<sub>3</sub> MOSFETs.

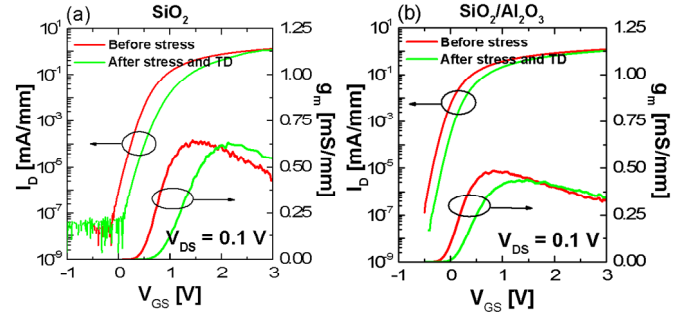


Figure 5. Room temperature subthreshold and transconductance characteristics before and after  $V_{GS, \text{stress}} = 15$  V stress followed by thermal detrapping (TD) of (a) SiO<sub>2</sub> and (b) SiO<sub>2</sub>/Al<sub>2</sub>O<sub>3</sub> GaN MOSFETs in the linear regime ( $V_{DS} = 0.1$  V).

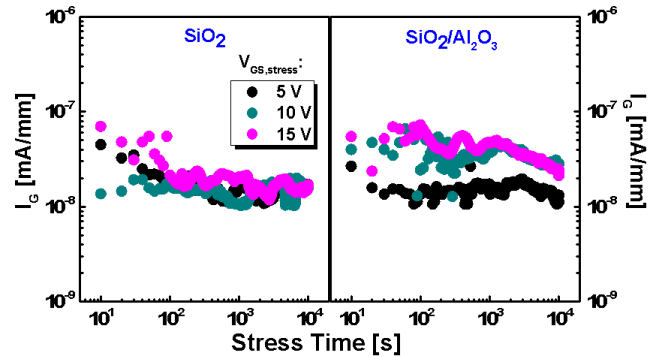


Figure 6. Stress time evolution of  $I_G$  for SiO<sub>2</sub> vs. SiO<sub>2</sub>/Al<sub>2</sub>O<sub>3</sub> GaN MOSFETs stressed at three different voltages at room temperature.

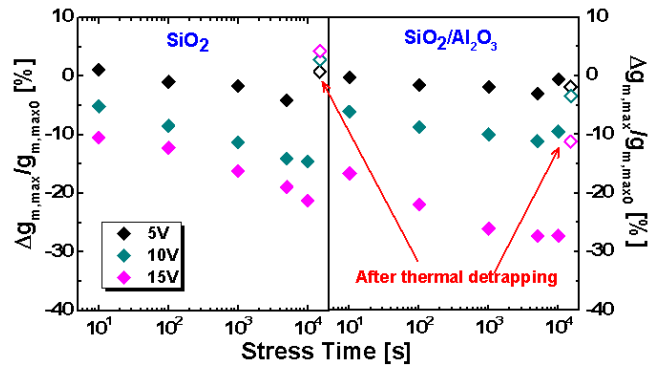


Figure 7. Stress time evolution of  $\Delta g_{m, \text{max}}/g_{m, \text{max}0}$  for SiO<sub>2</sub> vs. SiO<sub>2</sub>/Al<sub>2</sub>O<sub>3</sub> GaN MOSFETs stressed at three different voltages at room temperature. Data points are taken  $\sim 1$  sec after the stress is stopped. The last set of points (open symbols) are  $\Delta g_{m, \text{max}}/g_{m, \text{max}0}$  after thermal detrapping following a  $t_{\text{stress}} = 10,000$  s stress experiment.

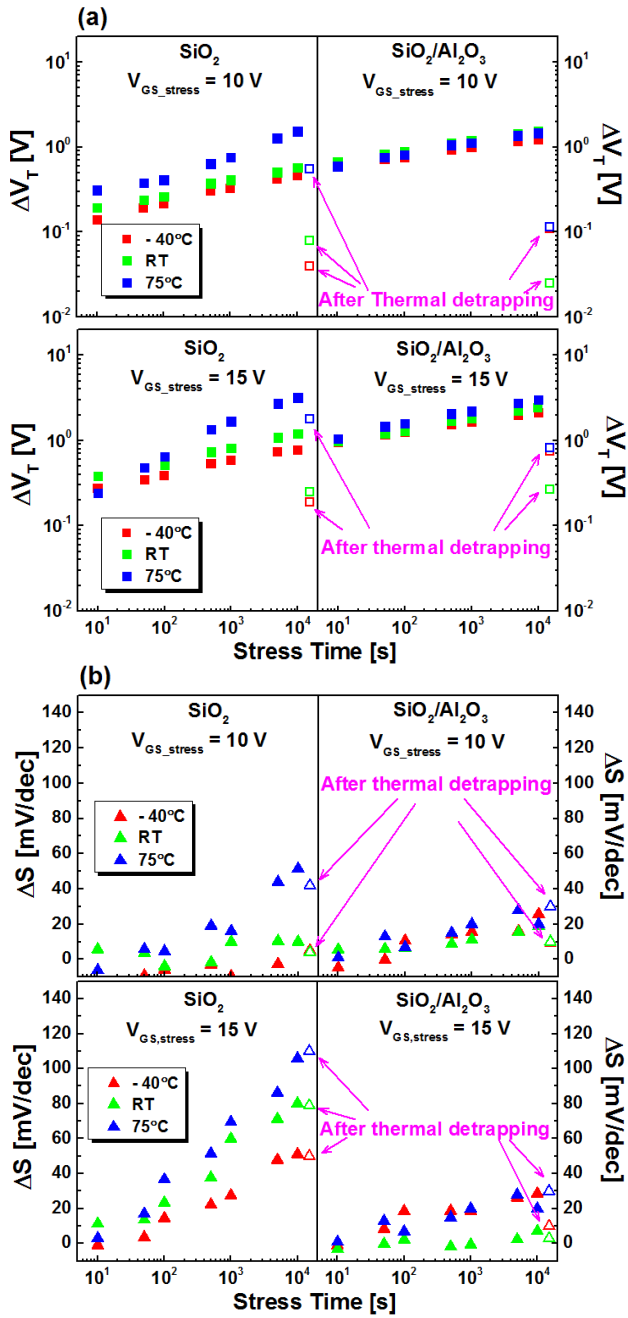


Figure 8. Stress time evolution of (a)  $\Delta V_T$  and (b)  $\Delta S$  for  $\text{SiO}_2$  vs  $\text{SiO}_2/\text{Al}_2\text{O}_3$  GaN MOSFETs at different  $T$ .  $V_{GS, \text{stress}} = 10$  and  $15$  V. Data points are taken 1 sec after the stress. The last set of points (open symbols) are  $\Delta V_T$  and  $\Delta S$  after a thermal detrapping step that follows stress experiments with  $t_{\text{stress}} = 10,000$  s.

### B. Impact of stress temperature

We have also studied the role of temperature during electrical stress. Fig. 8 shows  $\Delta V_T$  and  $\Delta S$  for  $V_{GS, \text{stress}} = 10$  and  $15$  V as a function of stress time at different temperatures. The stress and recovery phases take place at the same temperature. For  $\text{SiO}_2$ ,  $\Delta V_T$  and  $\Delta S$  increase with  $T$ , and the increase is more prominent at  $75^\circ\text{C}$ . On the other hand, the effect of  $T$  is small for  $\text{SiO}_2/\text{Al}_2\text{O}_3$  transistors at  $V_{GS, \text{stress}} = 10$  V, while at  $V_{GS, \text{stress}} = 15$  V, it becomes more prominent. However, this  $T$  dependence is less than in  $\text{SiO}_2$  transistors.

The recovery dynamics of  $V_T$  and  $S$  after  $V_{GS, \text{stress}} = 10$  V and  $15$  V at different  $T$  are graphed in Fig. 9. The recovery time constants of both  $\text{SiO}_2$  and  $\text{SiO}_2/\text{Al}_2\text{O}_3$  devices appear rather independent of temperature. The final values of  $\Delta V_T$  and  $\Delta S$  after thermal detrapping for  $t_{\text{stress}} = 10,000$  sec at various  $T$  are marked as open symbols (also summarized in Fig. 8 (open symbols)). Device characterization after thermal detrapping is made at the same temperature as the stress and recovery phases of the experiment.

For  $V_{GS, \text{stress}} = 10$  V, after thermal detrapping,  $\Delta V_T$  mostly recovers for the  $\text{SiO}_2/\text{Al}_2\text{O}_3$  transistors at all  $T$ . The  $\text{SiO}_2$  transistors recover for  $T \leq \text{RT}$  and only partially recover for  $T = 75^\circ\text{C}$ . Also at  $T = 75^\circ\text{C}$ ,  $S$  exhibits some permanent damage for  $\text{SiO}_2$  transistors. The  $\text{SiO}_2/\text{Al}_2\text{O}_3$  transistors, on the other hand, exhibit minor permanent damage in  $S$ .

For  $V_{GS, \text{stress}} = 15$  V, after thermal detrapping,  $\Delta V_T$  partially recovers for all transistors at all  $T$ . For the  $\text{SiO}_2$  transistors,  $\Delta V_T$  and  $\Delta S$  increase with  $T$ , and the increase is more prominent at  $75^\circ\text{C}$ . On the other hand, the effect of  $T$  is small for  $\text{SiO}_2/\text{Al}_2\text{O}_3$  transistors at  $V_{GS, \text{stress}} = 10$  V, while at  $V_{GS, \text{stress}} = 15$  V, it becomes more prominent. However, this  $T$  dependence is less than in  $\text{SiO}_2$  transistors.

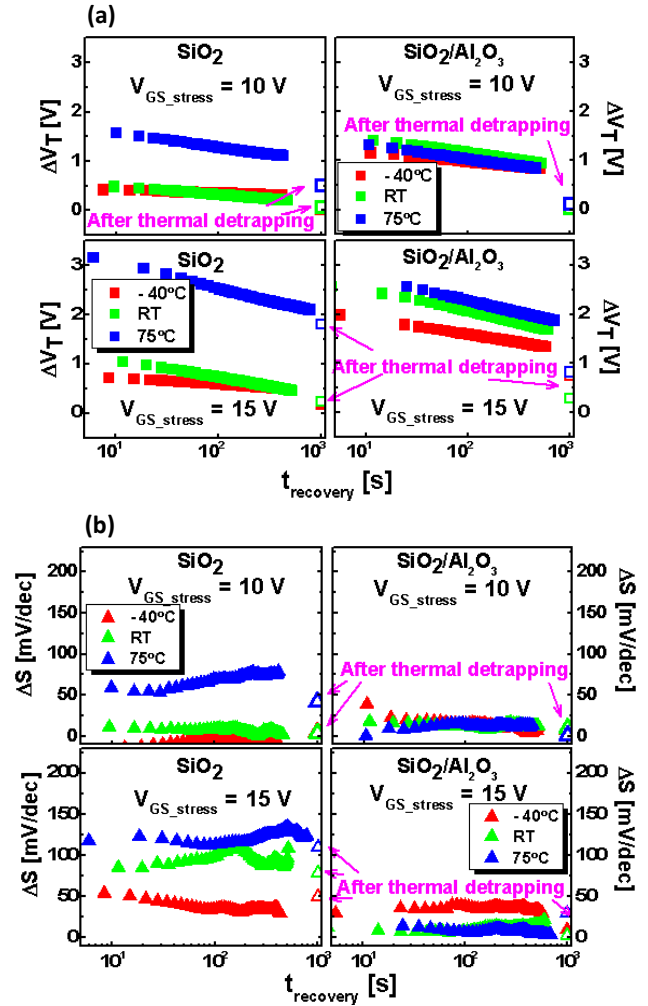


Figure 9. Recovery of (a)  $\Delta V_T$  and (b)  $\Delta S$  after a 10,000 sec stress at  $V_{GS, \text{stress}} = 10$  and  $15$  V for  $\text{SiO}_2$  vs.  $\text{SiO}_2/\text{Al}_2\text{O}_3$  GaN MOSFETs at  $-40^\circ\text{C}$ , RT and  $75^\circ\text{C}$ . The last set of points (open symbols) are  $\Delta V_T$  and  $\Delta S$  after subsequent thermal detrapping step.

there is more recovery for  $T \leq RT$  than for  $T = 75^\circ\text{C}$ . Also for  $\text{SiO}_2$  transistors,  $S$  exhibits permanent damage at all  $T$ , and the permanent damage increases with  $T$ . The  $\text{SiO}_2/\text{Al}_2\text{O}_3$  transistors again exhibit small permanent  $S$  damage. Again,  $I_G$  during stress tends to slightly decrease with stress time (not shown here).

#### IV. DISCUSSION

##### A. Mechanisms responsible for $V_T$ shift under PBTI

Our experiments allow us to postulate the mechanisms responsible for PBTI in GaN MOSFETs. Under benign stress ( $V_{GS, \text{stress}} \leq 10 \text{ V}$ ,  $T \leq RT$ ), the  $V_T$  shift is consistent with electron trapping in pre-existing oxide traps which is

characterized by a recoverable  $\Delta V_T$ . We conclude this from Figs. 8(a) and 9(a) that show  $\Delta V_T$  1 sec after different stress times. Within such a short time, little recovery takes place (Fig. 9 shows long recovery time constant). Stress time evolution of  $V_T$  follows a saturating log-time dependence, as observed in other MOS systems [15-17,19]. The slight decrease in  $I_G$  also indicates that no additional traps are generated in the oxide bulk [19].

Under harsher stress ( $V_{GS, \text{stress}} = 15 \text{ V}$ ), there is an additional permanent  $V_T$  shift for both gate dielectrics. Fig. 10 summarizes the final  $\Delta V_T$  and  $\Delta S$  after thermal detrapping for each stress segment at room temperature. Notably,  $\text{SiO}_2$  exhibits a significant permanent increase in  $S$  that suggests the generation of interface states. This is clearer in Fig. 11 that graphs the correlation between the final values of  $\Delta V_T$  and  $\Delta S$  after thermal detrapping following electrical stress at  $V_{GS, \text{stress}} = 15 \text{ V}$ . We see a strong linear correlation between the two. The proportionality constant seems to depend on  $T$ . Subthreshold swing degradation under PBTI has also been reported in  $\text{Al}_2\text{O}_3/\text{InGaAs}$  and  $\text{HfO}_2/\text{GaAs}$  MOSFETs [18,20]. In our composite oxide samples, the permanent degradation of  $S$  is minor and a clear correlation with  $\Delta V_T$  does not emerge.

An additional interesting observation in our work is that the evolution of  $g_{m, \text{max}}$  and  $V_T$  also seem correlated. This can be seen in Fig. 12 that graphs  $\Delta g_{m, \text{max}}/g_{m, \text{max}0}$  as a function of  $\Delta V_T$  after different stress periods for both dielectrics at room temperature. With the exception of long  $t_{\text{stress}}$  in  $\text{SiO}_2/\text{Al}_2\text{O}_3$  transistors, we observe a strong linear correlation between

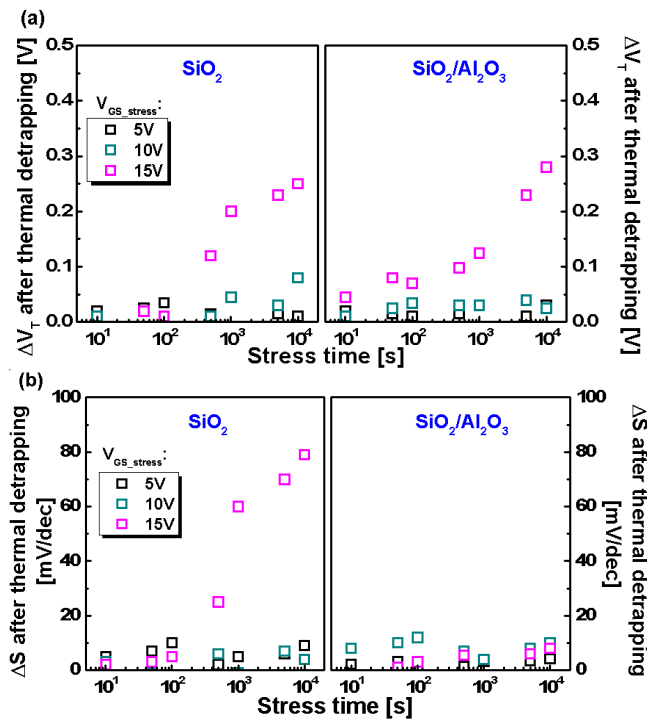


Figure 10. Final values of (a)  $\Delta V_T$  and (b)  $\Delta S$  after thermal detrapping vs. stress time for  $\text{SiO}_2$  vs.  $\text{SiO}_2/\text{Al}_2\text{O}_3$  GaN MOSFETs stressed at three different voltage at room temperature.

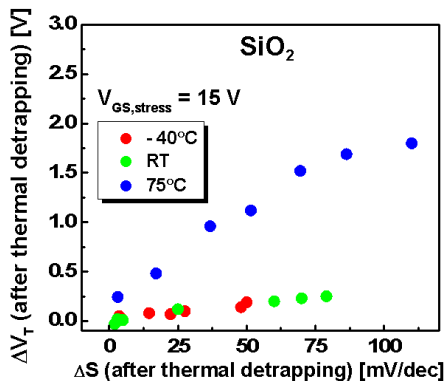


Figure 11. Evolution of final values of  $\Delta V_T$  vs.  $\Delta S$  for  $\text{SiO}_2$  transistors after  $V_{GS, \text{stress}} = 15 \text{ V}$  at different temperatures. All measurements were taken at room temperature after a thermal detrapping step.

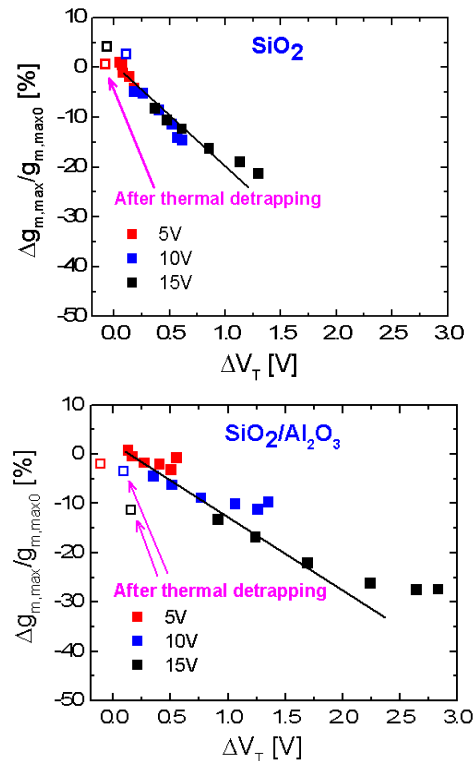


Figure 12. Evolution of  $(\Delta g_{m, \text{max}}/g_{m, \text{max}0})$  vs.  $\Delta V_T$  for (a)  $\text{SiO}_2$  and (b)  $\text{SiO}_2/\text{Al}_2\text{O}_3$  transistors during stress experiments of various durations at room temperature. The open symbols correspond to measurements after a thermal detrapping step following stress experiments with  $t_{\text{stress}} = 10,000 \text{ s}$ .

$\Delta g_{m,max}/g_{m,max0}$  and  $\Delta V_T$  at all  $V_{GS,stress}$  for both dielectrics. This suggests that  $\Delta V_T$  mostly originates in charge trapping in the oxide near the oxide/GaN interface or at the interface itself in interface states. Both are known to affect the mobility [17, 22]. For long stress times, trapping further away from the interface eventually takes place producing an additional  $V_T$  shift without a corresponding decrease in  $g_{m,max}$ .

$\Delta g_{m,max}/g_{m,max0}$  vs.  $\Delta V_T$  after thermal detrapping following stress experiments with  $t_{stress}=10,000$  s are also reported in Fig. 12 (open symbols). We see permanent degradation of  $\Delta g_{m,max}/g_{m,max0}$  for the  $\text{SiO}_2/\text{Al}_2\text{O}_3$  transistors after  $V_{GS,stress} = 15$  V, which suggests new trap generation in the oxide near the oxide/GaN interface that causes permanent mobility degradation [17, 23, 24]. Because these trap sites are far from the gate metal, no significant changes are observed in the evolution of  $I_G$ . This oxide trap generation close to the oxide/GaN interface also explains the incomplete recovery after thermal detrapping of  $V_T$  for the  $\text{SiO}_2/\text{Al}_2\text{O}_3$  transistor after  $V_{GS,stress} = 15$  V that is apparent in Figs. 3(a), 8 and 12.

### B. Model

The observations that we have made here are largely consistent with similar experiments on Si, SiC and III-V MOSFETs [15-18].  $\Delta V_T$  appears to take place due to two mechanisms. One component,  $\Delta V_{Tox}$ , arises from electron trapping in the oxide. A second component,  $\Delta V_{Tit}$ , which correlates to  $\Delta S$ , originates in interface state generation and population.

We postulate that, to the first order, for  $\text{SiO}_2$  MOSFETs,  $\Delta V_{Tit}$  is the permanent portion of  $\Delta V_T$  that is left after thermal detrapping. This hypothesis is supported by the linear dependence that we observe between the permanent portion of  $\Delta V_T$  and of  $\Delta S$  (Fig. 11) which strongly suggests the generation of interface states. Similar observation of interface state generation has been made in Si systems after radiation [25], and also in FinFETs [26] where the [110] Si surface has a much higher Si-H bond density compared to the [100] surface [27]. It is well known that hydrogen used to passivate dangling bonds at the oxide-semiconductor interface can escape leading to the creation of interface states [28]. The remainder of  $\Delta V_T$  for the  $\text{SiO}_2$  transistors is  $\Delta V_{Tox}$ .

For  $\text{SiO}_2/\text{Al}_2\text{O}_3$  MOSFETs,  $\Delta V_{Tox}$  dominates and the permanent portion of  $\Delta V_T$  is oxide trap generation close to the oxide/GaN interface, as observed in InGaAs n-MOSFETs [17]. The interface state generation in  $\text{SiO}_2/\text{Al}_2\text{O}_3$  appears negligible, which is consistent with literature on high-K dielectrics [19].

With this interpretation, we find that  $\Delta V_{Tox}$  in both transistor types follows a classic saturating log-time dependence, as observed in other MOS systems. The model of this power law relationship is described by [19]:

$$\Delta V_{Tox} = \Delta V_{max} \cdot \left\{ 1 - \exp \left( - \left( \frac{t_{stress}}{\tau_0} \right)^\beta \right) \right\} \quad (1)$$

where  $\Delta V_{Tox}$  is the  $V_T$  shift due to oxide trapping.  $\Delta V_{max}$  is a function of total trap density and the centroid of the trap-charge distribution in space,  $\beta$  describes the trap energy distribution, and  $\tau_0$  is the time constant of the traps. For the  $\text{SiO}_2/\text{Al}_2\text{O}_3$  system, an exponent  $\beta = 0.22 \sim 0.25$  and  $\tau_0 = 200$  s gives an excellent fit to the entire data set at all T (Fig. 13). For the  $\text{SiO}_2$  system,  $\beta = 0.25$  and  $\tau_0 = 150$  s provides a reasonable match at

all T. The  $\beta$  values extracted here are close to those reported in the literature:  $\beta = 0.2 \sim 0.32$  [15-18].

For  $V_{GS,stress} = 15$  V in  $\text{SiO}_2$  MOSFETs, there is an additional component of  $\Delta V_T$  that we attribute to interface state generation ( $\Delta V_{Tit}$ ) [25]. As we can see in Fig. 11, there is a linear dependence between  $\Delta V_{Tit}$  and  $\Delta S$  which is precisely what is expected from a simple model [25].

For  $\text{SiO}_2/\text{Al}_2\text{O}_3$  MOSFETs, there is also a non-recoverable  $\Delta V_T$  component for  $V_{GS,stress} = 15$  V. This permanent  $\Delta V_T$  is accompanied by a non-recoverable  $g_{m,max}$  degradation as shown in Fig. 7. This indicates permanent mobility degradation. Studies have attributed this to the generation of oxide traps close to the oxide/semiconductor interface [17,22].

## V. CONCLUSIONS

In conclusion, we have studied PBTI of GaN MOSFETs with  $\text{SiO}_2$  and  $\text{SiO}_2/\text{Al}_2\text{O}_3$  as gate dielectric. For  $\text{SiO}_2$  MOSFETs, a positive  $\Delta V_T$  is caused by a combination of electron trapping in pre-existing oxide traps and interface trap generation. For  $\text{SiO}_2/\text{Al}_2\text{O}_3$  MOSFETs, a positive  $V_T$  shift is

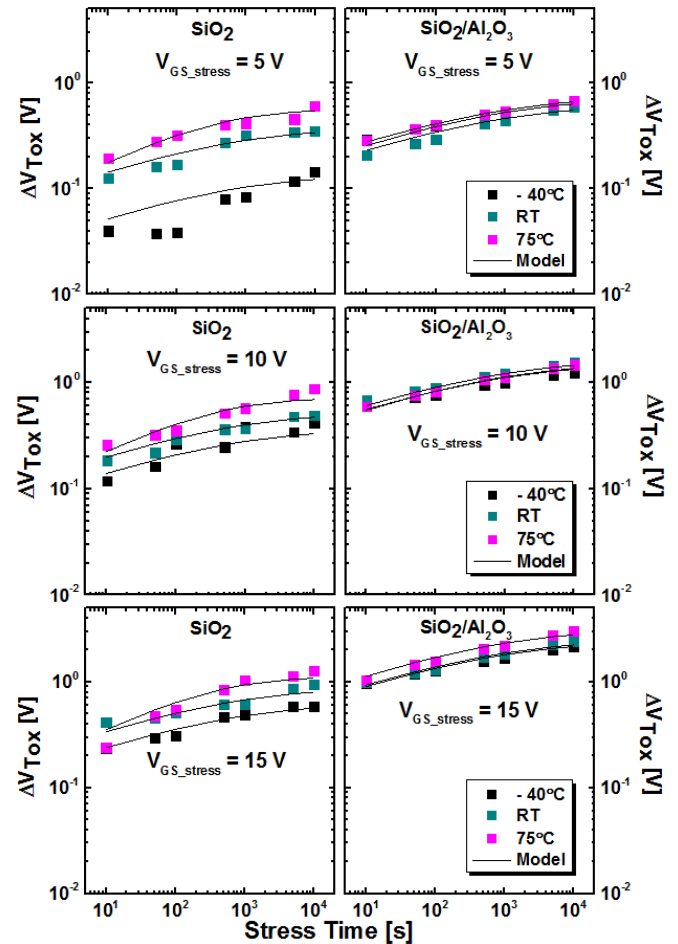


Figure 13. Model vs. experiments of  $\Delta V_{Tox}$  evolution with stress time, for  $V_{GS,stress} = 5$  V, 10 V, and 15 V, at different temperature.

due to trapping in the pre-existing oxide traps and oxide trap

generation near the oxide/GaN interface. Permanent damage after harsh stress is induced in both oxide systems but of a different nature. In SiO<sub>2</sub> transistors, non-recoverable interface state generation takes place. In SiO<sub>2</sub>/Al<sub>2</sub>O<sub>3</sub> MOSFETs, trap states in the oxide close to the semiconductor interface appear to be created. These findings are consistent with studies on other semiconductor material systems.

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